

# Points fixes et centres de libration *dans les chaînes de résonance*

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## Exoplanetary systems often are in resonance chains

- ▶ Examples: Kepler-223 (3 : 4 : 6 : 8) and TOI-178 (0 : 2 : 4 : 6 : 9 : 12).
- ▶ In the solar system: The four galilean satellites (1 : 2 : 4 : 0).
- ▶ A generic  $n$ -planets chain is denoted  $p_1 : p_2 : \dots : p_{n-1} : p_n$ .

## Definition

- ▶ A pair  $(i, j)$  of planet is resonant if  $T_i/T_j \sim p_i/p_j \in \mathbb{Q}$ .
- ▶ Or equivalently (Kepler third law) if  $a_i/a_j \sim (p_i/p_j)^{2/3}$ .

## Two planets in resonance $p : p + 1$ (planar case)

$$\boxed{\frac{T_2}{T_1} \approx \frac{p+1}{p}}$$

e.g.  $T_2/T_1 \approx 3/2$  for Kepler-1972

► We use Poincaré's canonical variables

$$\Lambda_j = \beta_j \sqrt{\mu_j a_j} \quad \leftrightarrow \quad \lambda = \text{mean longitude},$$

$$D_j = \Lambda_j \left(1 - \sqrt{1 - e_j^2}\right) \quad \leftrightarrow \quad -\varpi_j = -\text{longitude of the periapsis},$$

where  $\beta_j = m_0 m_j / (m_0 + m_j) \approx m_j$  and  $\mu_j = \mathcal{G} \sqrt{m_0 + m_j}$ .

## Hamiltonian

$$\mathcal{H} = \mathcal{H}_K(\Lambda_1, \Lambda_2) + \mathcal{H}_P(\Lambda_1, \Lambda_2, D_1, D_2; \lambda_1, \lambda_2, -\varpi_1, -\varpi_2),$$

$$\mathcal{H}_K(\Lambda_1, \Lambda_2) = - \sum_{j=1}^2 \frac{\beta_j^3 \mu_j^2}{2\Lambda_j^2}.$$

- The perturbation is expanded in a Fourier serie of eccentricities

$$\mathcal{H}_P = \sum_{\mathbf{k} \in \mathbb{Z}^2} \left( \sum_{\mathbf{q} \in \mathbb{N}^4} \Xi(a_1/a_2) X_1^{q_1} X_2^{q_2} \bar{X}_1^{q_3} \bar{X}_2^{q_4} \right) e^{i(k_1 \lambda_1 + k_2 \lambda_2)},$$

where  $X_j = e_j \exp i\varpi_j + \mathcal{O}(e_j^3)$ .

**Invariance by rotation** (D'Alembert rule):

$$k_1 + k_2 + q_1 + q_2 - q_3 - q_4 = 0$$

## Averaging and truncation

- ▶ Most of the angles  $k_1\lambda_1 + k_2\lambda_2$  are fast-circulating.
- ▶ We only keep the slow ones
- ▶ Example: Resonance  $2 : 3 \rightarrow$  We keep  $2\lambda_1 - 3\lambda_2$  and the multiples.
- ▶ Order  $r$  in eccentricity: We only keep  $|q_1| + |q_2| + |q_3| + |q_4| \leq r$

## First order in eccentricity in resonance $p : p + 1$

$$\mathcal{H}_P = \frac{Gm_1m_2}{a_2} \left[ f_1 \sqrt{\frac{2D_1}{\Lambda_1}} \cos(p\lambda_1 - (p+1)\lambda_2 + \varpi_1) + f_2 \sqrt{\frac{2D_2}{\Lambda_2}} \cos(p\lambda_1 - (p+1)\lambda_2 + \varpi_2) \right].$$

## Reduction to one degree of freedom

- ▶ Two angles of interest  $\sigma_j = p\lambda_1 - (p+1)\lambda_2 + \varpi_j$
- ▶ Transformation  $(e_1, e_2; \sigma_1, \sigma_2) \rightarrow (R, S; r, s)$
- ▶  $s$  absent from the Hamiltonian  $\rightarrow S$  is a conserved quantity.

## Conserved quantities

- ▶ Total angular momentum  $G = m_1 \sqrt{\mu a_1 (1 - e_1^2)} + m_2 \sqrt{\mu a_2 (1 - e_2^2)}$
- ▶ Scaling parameter  $\Gamma = (p+1) m_1 \sqrt{\mu a_1} + p m_2 \sqrt{\mu a_2}$
- ▶  $S \propto \Upsilon = \frac{e_{1,\text{eq}}}{e_{2,\text{eq}}} e_2^2 + \frac{e_{2,\text{eq}}}{e_{1,\text{eq}}} e_1^2 + 2e_1 e_2 \cos(\varpi_1 - \varpi_2)$

$$\text{e.g. } \Upsilon_{2:3} = 1.07148 \frac{m_1}{m_2} e_1^2 + 0.93329 \frac{m_2}{m_1} e_2^2 + 2e_1 e_2 \cos(\varpi_1 - \varpi_2)$$

## Reduction to one parameter

$$\mathcal{H} = \alpha R - \beta R^2 + \gamma \sqrt{2R} \cos r$$

$$\beta(p, m_1, m_2), \quad \gamma(p, m_1, m_2), \quad \alpha(p, m_1, m_2, G + S)$$

- Rescaling action and time :  $\Sigma = (2\beta/\gamma)^{2/3} R$  and  $\tau = (\beta\gamma^2/4)^{1/3} t$

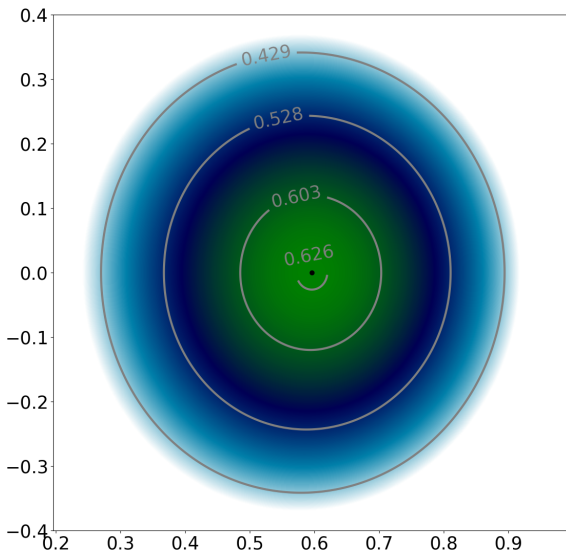
$$\boxed{\mathcal{H} = 3\delta\Sigma - \Sigma^2 + 2\sqrt{2\Sigma} \cos \sigma} \quad \delta = \alpha \left( \frac{4}{27\beta\gamma^2} \right)^{1/3}$$

- Using cartesian coordinates  $X = \sqrt{2\Sigma} \cos \sigma$  and  $Y = \sqrt{2\Sigma} \sin \sigma$

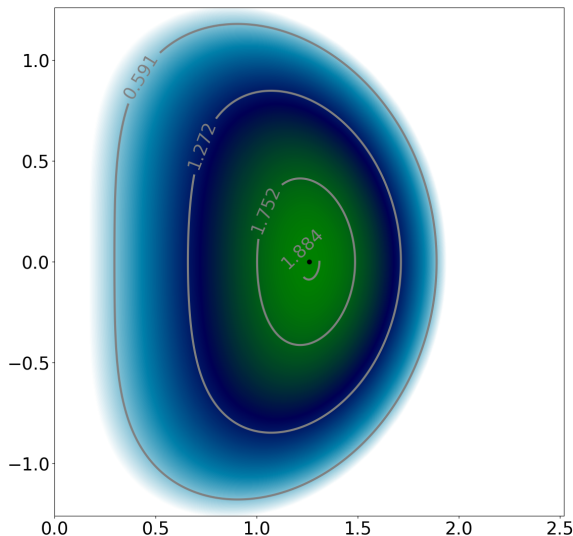
$$\boxed{\mathcal{H} = 2X - \frac{1}{4} (X^2 + Y^2 - 3\delta)^2}$$

- $|X| \sim \mathcal{O} \left( e_j \left( \frac{m_0}{m_j} \right)^{1/3} \right)$  and  $|Y| \sim \mathcal{O} \left( e_j \left( \frac{m_0}{m_j} \right)^{1/3} \right)$

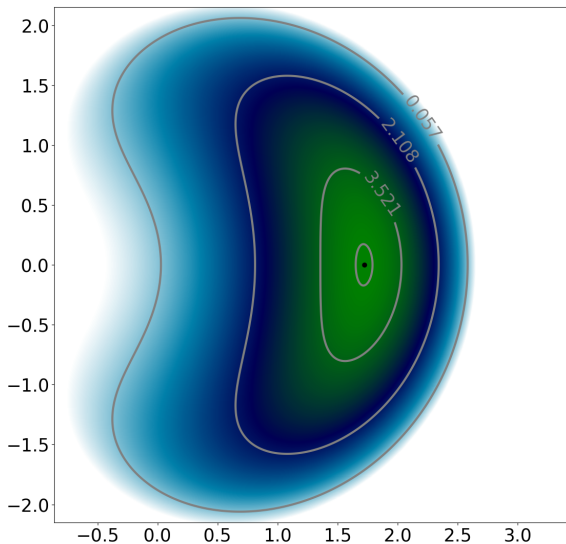
$$\delta = -1$$



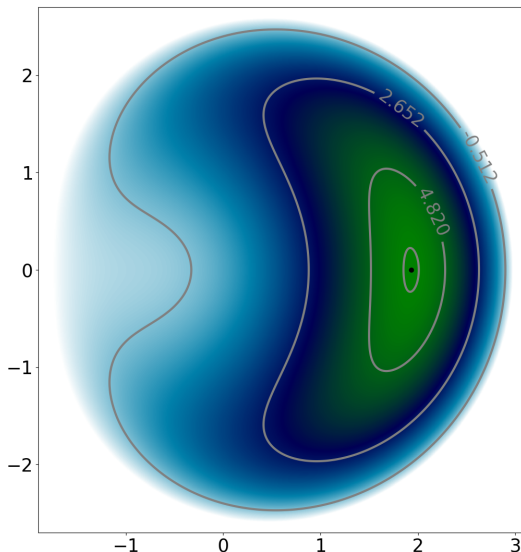
$$\delta = 0$$



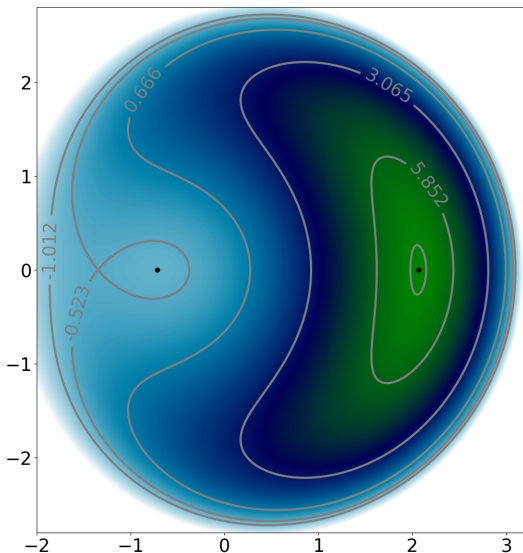
$$\delta = 0.6$$



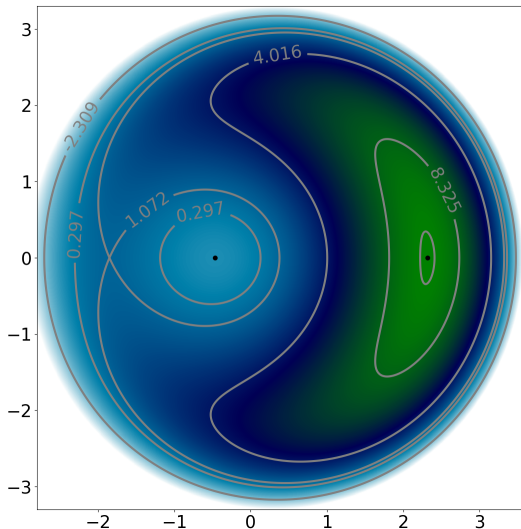
$$\delta = 0.9$$



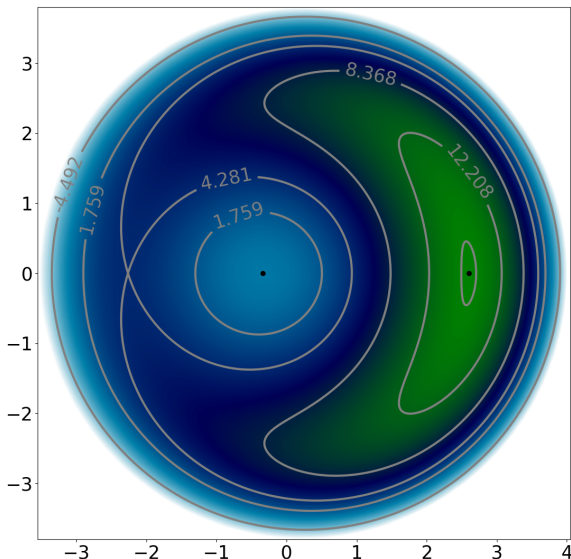
$$\delta = 1.1$$



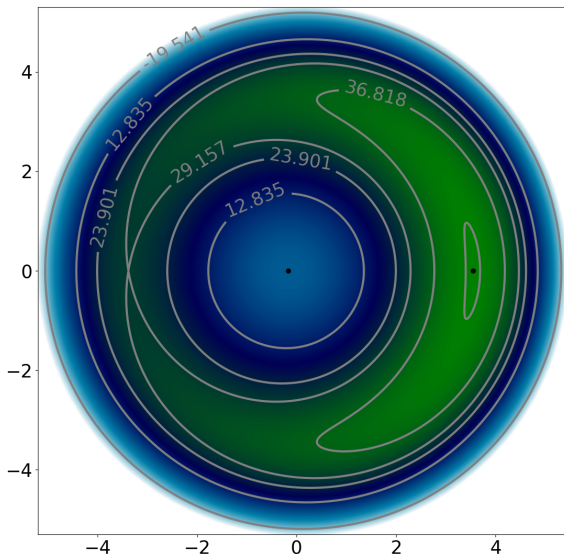
$$\delta = 1.5$$



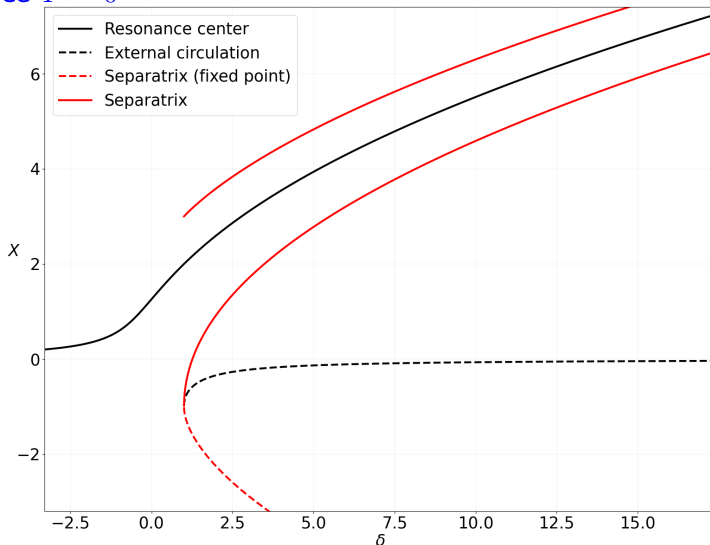
$$\delta = 2$$



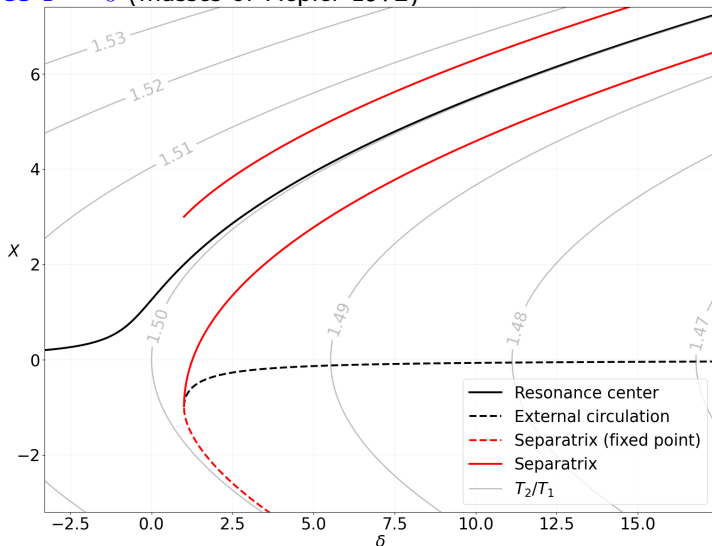
$$\delta = 4$$



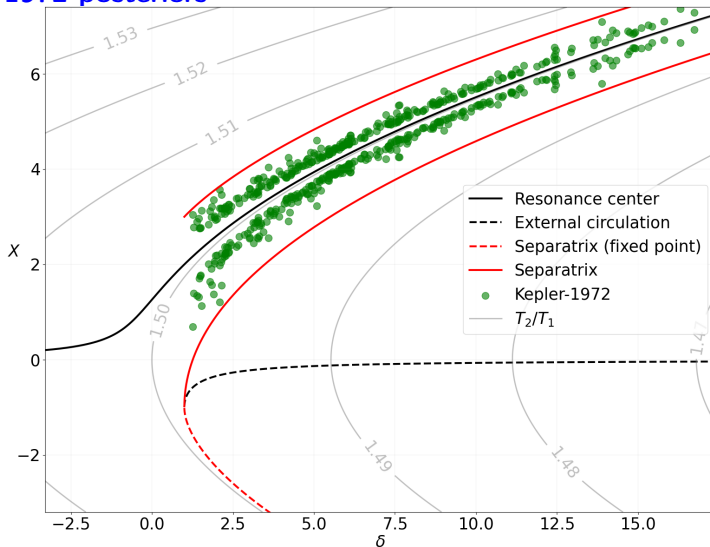
## Subspace $Y = 0$



## Subspace $Y = 0$ (masses of Kepler-1972)



## Kepler-1972 posteriors



## Overly simplified model. What about a real system?

Many approximations, including:

- ▶ Perturbations of the pair by other planets when  $N > 2$  are disregarded.
- ▶ Model limited to first-order in eccentricity
- ▶ Perturbative part simply evaluated at  $a_1/a_2 = (p/(p+1))^{2/3}$
- ▶ Averaging over the fast angles of the Hamiltonian
- ▶ Contributions in  $\mathcal{O}(m_j/m_0)^2$  disregarded

## I am creating a software that

- ▶ Automatically studies a given chain  $p_1 : p_2 : \dots : p_n$
- ▶ Can manipulate both the averaged and complete Hamiltonian
- ▶ Can find fixed points
- ▶ Can find *libration centers*

## Aptidal (name might change, IDK)

The resonance chain to be studied is 3:4:6  
The system currently has 6 degrees of freedom.

Aptidal will perform a canonical transformation  $(\text{lbd}, -\text{vrp}; \text{Lbd}, D) \rightarrow (\text{phi}, \text{sig}; \text{Phi}, D)$   
where  $(\text{lbd}, -\text{vrp}; \text{Lbd}, D)$  are the traditional Poincaré variables.  
 $\text{lbd}$  ( $\lambda$ ) is the mean longitude,  $\text{vrp}$  ( $\varpi$ ) is the longitude of the ascending node,  $\text{Lbd}$  ( $\text{resp. Phi}$ )  
is the action conjugated to  $\lambda$  ( $\text{resp. phi}$ ) and  $D$  (the AMD) is conjugated to both  $\sigma$  and  $-\varpi$ .

The transformation reads

$$\begin{aligned}\text{phi}_1 &= \text{lbd}_1 - 2 \text{lbd}_2 + \text{lbd}_3 \\ \text{phi}_2 &= \text{lbd}_2 - \text{lbd}_3 \\ \text{phi}_3 &= -2 \text{lbd}_2 + 3 \text{lbd}_3 \\ \text{sig}_j &= -2 \text{lbd}_2 + 3 \text{lbd}_3 - \text{vrp}_j, \quad 1 \leq j \leq 3\end{aligned}$$

It is canonical if the actions are transformed according to (the  $D_j$  are unchanged)

$$\begin{aligned}\text{Phi}_1 &= \text{Lbd}_1 \\ \text{Phi}_2 &= 4 \text{Lbd}_1 + 3 \text{Lbd}_2 + 2 \text{Lbd}_3 \\ \text{Phi}_3 &= \text{Lbd}_1 + \text{Lbd}_2 + \text{Lbd}_3 - D_1 - D_2 - D_3\end{aligned}$$

After averaging over the fast angle, the 4 remaining degrees of freedom are  
 $(\text{phi}_1; \text{Phi}_1), (\text{sig}_1; D_1), (\text{sig}_2; D_2), (\text{sig}_3; D_3)$

# Aptidal (name might change, IDK)

```

H = - sum_{1 <= j <= 3} beta_j**3 * mu_j**2 / (2*Lbd_j**2)

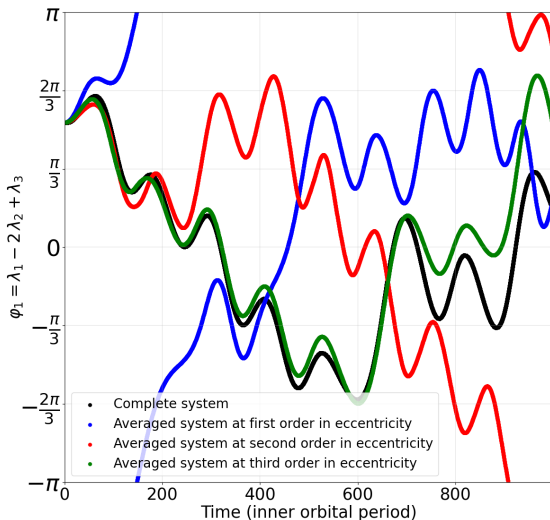
+ G*m_1*m_2/a_2 * (
  - 2.1973798573 * sqrt(2*D_1/Lbd_1)**2
  - 2.1973798573 * sqrt(2*D_2/Lbd_2)**2
  + 4.0222128577 * sqrt(2*D_1/Lbd_1) * sqrt(2*D_2/Lbd_2) * cos(vrp_1 - vrp_2)
  + 2.7760167116 * sqrt(2*D_1/Lbd_1) * cos(3 lbd_1 - 4 lbd_2 + vrp_1)
  - 3.2201137810 * sqrt(2*D_2/Lbd_2) * cos(3 lbd_1 - 4 lbd_2 + vrp_2)
  - 10.472815452 * sqrt(2*D_1/Lbd_1)**2 * cos(6 lbd_1 - 8 lbd_2 + 2 vrp_1)
  + 24.574405534 * sqrt(2*D_1/Lbd_1) * sqrt(2*D_2/Lbd_2) * cos(6 lbd_1 - 8 lbd_2 + vrp_1 + vrp_2)
  - 14.385904679 * sqrt(2*D_2/Lbd_2)**2 * cos(6 lbd_1 - 8 lbd_2 + 2 vrp_2)
)

+ G*m_1*m_3/a_3 * (
  - 0.3693855360 * sqrt(2*D_1/Lbd_1)**2
  - 0.3693855360 * sqrt(2*D_3/Lbd_3)**2
  + 0.5435235753 * sqrt(2*D_1/Lbd_1) * sqrt(2*D_3/Lbd_3) * cos(vrp_1 - vrp_3)
  + 1.1556528318 * sqrt(2*D_1/Lbd_1) * cos(lbd_1 - 2 lbd_3 + vrp_1)
  - 0.3896164810 * sqrt(2*D_3/Lbd_3) * cos(lbd_1 - 2 lbd_3 + vrp_3)
  - 1.6079268247 * sqrt(2*D_1/Lbd_1)**2 * cos(2 lbd_1 - 4 lbd_3 + 2 vrp_1)
  + 4.7599824848 * sqrt(2*D_1/Lbd_1) * sqrt(2*D_3/Lbd_3) * cos(2 lbd_1 - 4 lbd_3 + vrp_1 + vrp_3)
  - 3.4807825350 * sqrt(2*D_3/Lbd_3)**2 * cos(2 lbd_1 - 4 lbd_3 + 2 vrp_3)
)

+ G*m_2*m_3/a_3 * (
  - 1.0937567733 * sqrt(2*D_2/Lbd_2)**2
  - 1.0937567733 * sqrt(2*D_3/Lbd_3)**2
  + 1.8878074721 * sqrt(2*D_2/Lbd_2) * sqrt(2*D_3/Lbd_3) * cos(vrp_2 - vrp_3)
  + 1.9595502701 * sqrt(2*D_2/Lbd_2) * cos(2 lbd_2 - 3 lbd_3 + vrp_2)
  - 2.4205566721 * sqrt(2*D_3/Lbd_3) * cos(2 lbd_2 - 3 lbd_3 + vrp_3)
  - 5.0371064812 * sqrt(2*D_2/Lbd_2)**2 * cos(4 lbd_2 - 6 lbd_3 + 2 vrp_2)
  + 12.655477566 * sqrt(2*D_2/Lbd_2) * sqrt(2*D_3/Lbd_3) * cos(4 lbd_2 - 6 lbd_3 + vrp_2 + vrp_3)
  - 7.9162336115 * sqrt(2*D_3/Lbd_3)**2 * cos(4 lbd_2 - 6 lbd_3 + 2 vrp_3)
)

```

## Can integrate the model and the real system automatically



## Fixed points vs libration center ? Chain 3 : 4 : 6 : 8 Fixed points

- ▶  $\phi_1 = \lambda_1 - 2\lambda_2 + \lambda_3 = \text{Cte}$
- ▶  $\phi_2 = \lambda_2 - 3\lambda_3 + 2\lambda_4 = \text{Cte}$
- ▶  $\phi_3 = \lambda_3 - \lambda_4$
- ▶  $\phi_4 = -3\lambda_3 + 4\lambda_4$
- ▶  $\sigma_j = -3\lambda_3 + 4\lambda_4 - \varpi_j = \text{Cte}$

### Libration center

- ▶  $\phi_1 = \lambda_1 - 2\lambda_2 + \lambda_3 =$  is periodic of period  $\nu_3$
- ▶  $\phi_2 = \lambda_2 - 3\lambda_3 + 2\lambda_4 =$  is periodic of period  $\nu_3$
- ▶  $\phi_3 = \lambda_3 - \lambda_4 = \nu_3 t +$  periodic of period  $\nu_3$
- ▶  $\phi_4 = -3\lambda_3 + 4\lambda_4 = \nu_4 t +$  periodic of period  $\nu_3$
- ▶  $\sigma_j = -3\lambda_3 + 4\lambda_4 - \varpi_j =$  is periodic of period  $\nu_3$

## Can find fixed points of the model

Starting the search for a fixed point.

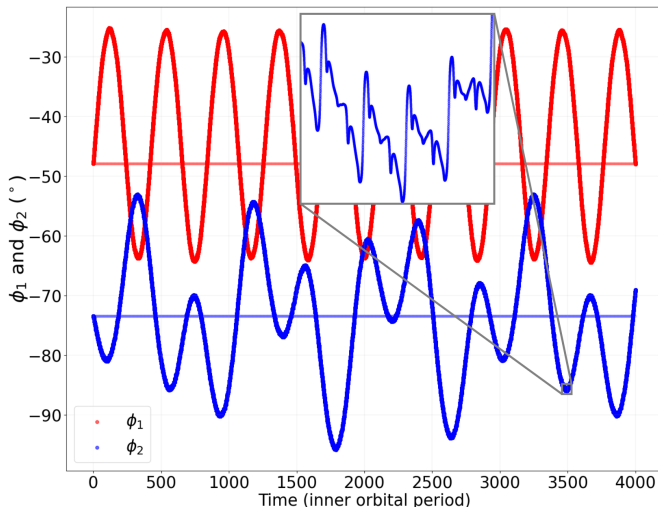
```
Iteration n° 0 : (lbd_1, lbd_2, lbd_3) = (0.9999991553361, 0.00000000000000, 0.00000000000000) x pi
                (vrp_1, vrp_2, vrp_3) = (0.6816892691523, 0.00000000000000, 0.9999991553361) x pi
                (a_1, a_2, a_3)      = (0.9999646060000, 1.2157125170000, 1.6045573520000)
                (e_1, e_2, e_3)      = (0.0039254030000, 0.0035638750000, 0.0014761090000)
```

```
Iteration n° 1 : (lbd_1, lbd_2, lbd_3) = (1.0000048947343, 0.00000000000000, 0.00000000000000) x pi
                (vrp_1, vrp_2, vrp_3) = (0.9999018601172, 0.0000772588654, 0.9998723307605) x pi
                (a_1, a_2, a_3)      = (0.9999528690489, 1.2157312263377, 1.6045437220038)
                (e_1, e_2, e_3)      = (0.0039117085973, 0.0035517536071, 0.0014708166399)
```

```
Iteration n° 6 : (lbd_1, lbd_2, lbd_3) = (0.9999999999989, 0.00000000000000, 0.00000000000000) x pi
                (vrp_1, vrp_2, vrp_3) = (-0.9999999999973, 0.00000000000021, 0.9999999999999) x pi
                (a_1, a_2, a_3)      = (0.9999646061216, 1.2157125167514, 1.6045573522558)
                (e_1, e_2, e_3)      = (0.0039254046462, 0.0035638770699, 0.0014761094377)
```

## Finding a fixed point

- By successive average of a variable  $f(t)$  of the Hamiltonian





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## Converging to a libration center

- ▶ The system is integrated in the complete system using a fixed point of the model as initial condition
- ▶ 
$$f(t) = \sum_{\mathbf{k} \in \mathbb{Z}^8} f_{\mathbf{k}} e^{i\mathbf{k} \cdot \boldsymbol{\omega} t}$$
- ▶  $\boldsymbol{\omega} = (\nu_1, \nu_2, \dots, \nu_8)$  vector of fundamental frequencies.
- ▶  $f_{\mathbf{k}}$  is set to 0 except if  $\mathbf{k} = (0, 0, k, 0, 0, 0, 0, 0)$  for  $k \in \mathbb{Z}$ .
- ▶ Done on all variables of the Hamiltonian
- ▶ Defines a new initial condition from which the process is reiterated.

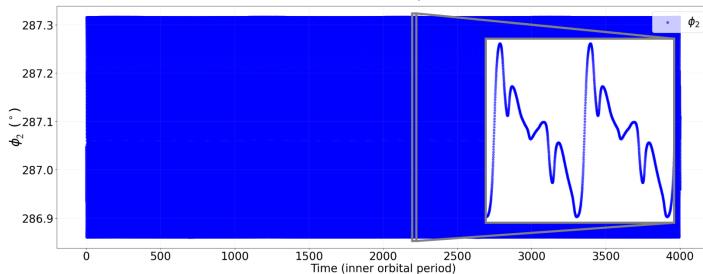
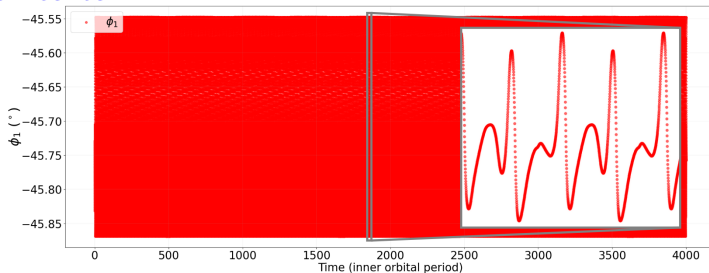


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## Computing $f_{\mathbf{k}}$ : NAFF algorithm

- ▶ The system is integrated from  $t = -T$  to  $t = T$
- ▶  $\nu = \mathbf{k} \cdot \boldsymbol{\omega}$
- ▶  $f_{\mathbf{k}} = \frac{1}{2T} \int_{-T}^T f(t) e^{-i\nu t} H_p(t) dt$
- ▶  $\nu = k\nu_3 = (0, 0, k, 0, 0, 0, 0, 0) \cdot (\nu_1, \nu_2, \nu_3, \nu_4, \nu_5, \nu_6, \nu_7, \nu_8)$
- ▶  $\nu_3$  obtained from a linear fit of  $\phi_3(t) = \nu_3 t + \phi_3(0)$





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## Periodic orbit finding

- ▶ Variables of the Hamiltonian periodic with frequency  $\nu_3 = 0.79780724$
- ▶ In a frame rotation with the pericenters at frequency  $\nu_4 = 0.09929880$
- ▶  $\nu_3/\nu_4 = 8.0344$
- ▶ If  $\nu_3/\nu_4 \in \mathbb{Q}$ , then the motion is periodic
- ▶ A whole family of libration centers parameterized by the total angular momentum
- ▶ Periodic orbits can be found by varying the angular momentum until  $\nu_3/\nu_4 \in \mathbb{Q}$